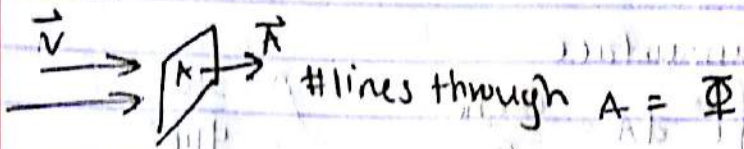
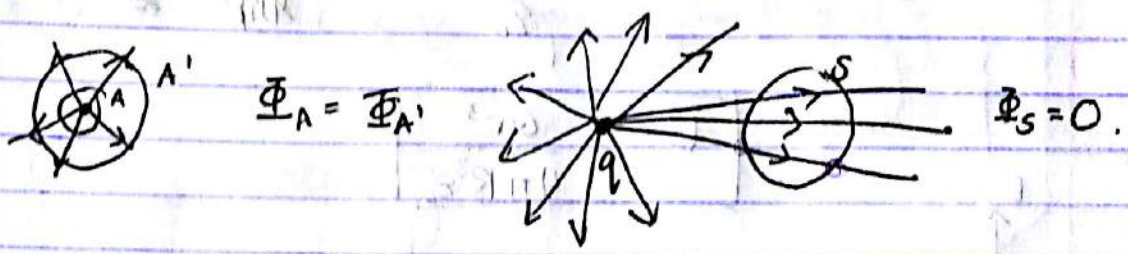


Review: Flux

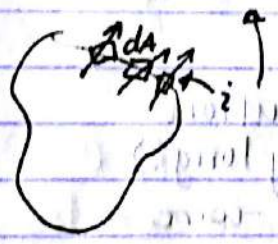


$\Phi = \vec{E} \cdot \vec{A}$, but there is a dependence on θ (between \vec{E} & \vec{A})
 $\hookrightarrow \vec{E} \cdot \vec{A} = \Phi$



$\Phi_E = \frac{Q_{\text{enclosed}}}{\epsilon_0} \Rightarrow \oint \vec{E} \cdot d\vec{A} = \frac{Q_{\text{enclosed}}}{\epsilon_0}$ (uniform charge distribution)

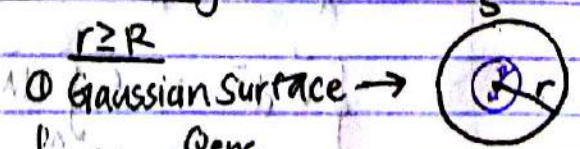
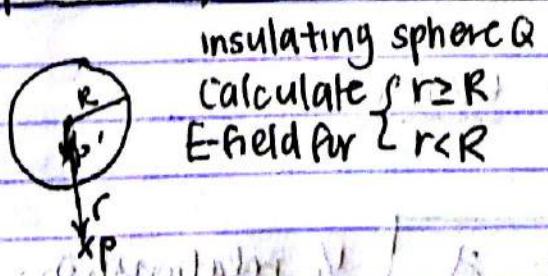
$\Phi_{\text{tot}} = \sum_i \Phi_i = \oint \vec{E} \cdot d\vec{A}$



Non uniform charge distribution

$\Phi_E = \oint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} \int_V \rho dV$ (charge distribution)

spherical symmetric distribution of charge



$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{\text{enc}}}{\epsilon_0}$

$E \oint dA = \frac{Q}{\epsilon_0}$

$E = \frac{Q}{4\pi\epsilon_0 r^2}$

like a point charge

$E(4\pi r^2) = Q/\epsilon_0$

$r < R$



Gaussian surface

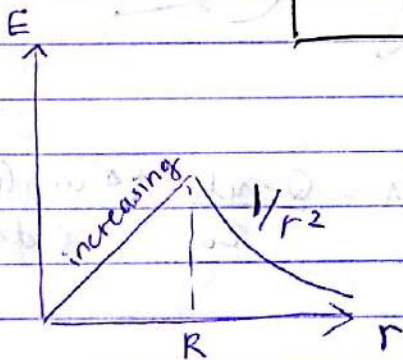
$$\Phi_E = \oint_{S'} \vec{E} \cdot d\vec{A} = \frac{Q_{enc}}{\epsilon_0}$$

$$\vec{E} 4\pi r^2 = \frac{\rho V_{S'}}{\epsilon_0} \rightarrow$$

$$\rho = \frac{3Q}{4\pi R^3}$$

$$\frac{3Q}{4\pi R^3} \cdot \frac{4}{3}\pi r^3$$

$$\boxed{\vec{E} = \frac{Qr}{4\pi R^3 \epsilon_0}}$$



Cylindrical symmetrical charge distribution (very long!)



(very long!)
-wire

calculate \vec{E}

Gaussian surface

$$Q = \lambda l$$

charge density per unit length

$$\Phi_E = \Phi_A + \Phi_{A'}$$

$A' \cdot \vec{E}$ is 0!

$$= \oint_A \vec{E} \cdot d\vec{A} + \oint_{A'} \vec{E} \cdot d\vec{A} \rightarrow 0$$

$$\frac{Q_{enc}}{\epsilon_0} = \frac{\lambda l}{\epsilon_0}$$

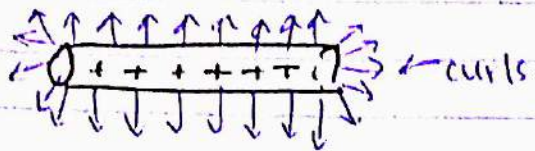
$$= \oint_A \vec{E} \cdot d\vec{A} = \vec{E} \int_A dA = E 2\pi r l$$

$$E 2\pi r l = \frac{\lambda l}{\epsilon_0}$$

$$\boxed{\vec{E} = \frac{\lambda}{2\pi r \epsilon_0}}$$

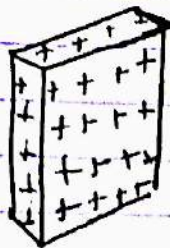
1/r relationship.

What if the wire is not infinite?



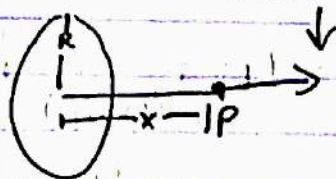
-when Gauss's Law fails, use Coulomb's Law!

\vec{E} for Plane Charge (non-conducting)



σ - charge per unit area distributed through V
Calculate E

Recall... Coulomb's:

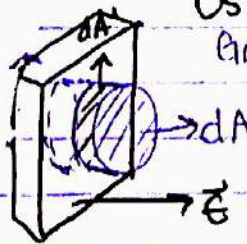


$$E_{\text{disk}} = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{x}{(x^2 + R^2)^{1/2}} \right)$$

$R \rightarrow \infty$, like a infinite plane $= \frac{\sigma}{2\epsilon_0}$

Using Gauss's Law:

Gaussian surface - cylinder



$$\Phi_{\text{TOT}} = \Phi_{A_{\text{TOT}}} + \Phi_{A'_{\text{TOT}}}$$

$$= \int_{A_{\text{TOT}}} \vec{E} \cdot d\vec{A} = \vec{E} \int_{A_{\text{TOT}}} d\vec{A} = 2\vec{E} \int_A d\vec{A} = 2EA$$

$$\frac{Q_{\text{enc}}}{\epsilon_0} = \frac{\sigma A}{\epsilon_0} = 2EA$$

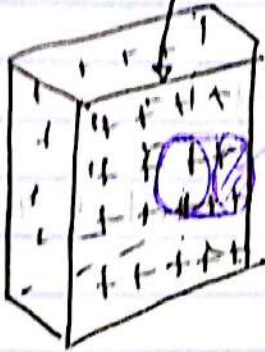
$$E = \frac{\sigma}{2\epsilon_0}$$

\vec{E} for conducting plane



σ

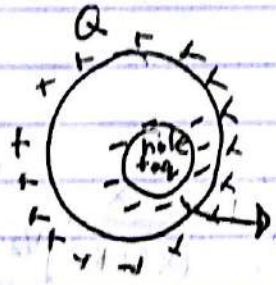
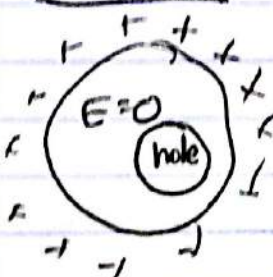
$\vec{E} = 0$ inside



$$\Phi_G = EA = \frac{\sigma A}{\epsilon_0}$$

$$\vec{E} = \frac{\sigma}{\epsilon_0}$$

Conductor



$\times P E = ?$
 $\sim \frac{(Q+q)}{\epsilon_0}$
 surface encloses $Q \neq q$.

$\vec{E}_{inside} = ?$
 has to be 0.

conductors are lazy.
 They don't like fields.

charge will distribute itself so $\vec{E}_{inside} = 0$.
 (if not evenly dist, $F = qE$, so charges move.)