

## 16.9 Divergence Theorem

Flashback:

16.5, used Green's Theorem to derive:

$$\int_C \vec{F} \cdot d\vec{r} = \iint_D \text{div } \vec{F} \, dA.$$

Overview:

The divergence theorem is a higher-dimensional analog of Green's Theorem. Recall how solid regions of type 1, 2 and 3 in 15.7 are defined:

Type 1:  $E = \{(x, y, z) \mid (x, y) \in D, z_1(x, y) \leq z \leq z_2(x, y)\}$   
 Type 2: "  $(y, z) \in D, x_1(y, z) \leq x \leq x_2(y, z)\}$   
 Type 3: "  $(z, x) \in D, y_1(x, z) \leq y \leq y_2(x, z)\}$

Solids simultaneously of types 1, 2, and 3 are simple.

The Divergence theorem:

Let  $E$  be a simple sided solid region w/ boundary surface  $S$  and has positive orientation.

Let  $\vec{F}$  be a vector field on  $\mathbb{R}^3$  whose components have cont. partial deriv. on an open region containing  $E$ .

Then,

$$\iint_S \vec{F} \cdot d\vec{S} = \iiint_E \text{div } \vec{F} \, dV.$$

flux of  $\vec{F}$  across

$$S = dE$$

Proof:

If  $\vec{F} = P\hat{i} + Q\hat{j} + R\hat{k}$ , then

$$\text{div } \vec{F} = \frac{\partial P}{\partial x} + \frac{\partial Q}{\partial y} + \frac{\partial R}{\partial z}$$

$$\iiint_E \text{div } \vec{F} \, dV = \iiint_E \frac{\partial P}{\partial x} \, dV + \iiint_E \frac{\partial Q}{\partial y} \, dV + \iiint_E \frac{\partial R}{\partial z} \, dV.$$

On the other hand,

$$\begin{aligned} \iint_S \vec{F} \cdot d\vec{S} &= \iint_S \vec{F} \cdot \hat{n} \, dS = \iint_S (P\hat{i} + Q\hat{j} + R\hat{k}) \cdot \hat{n} \, dS \\ &= \iint_S P\hat{i} \cdot \hat{n} \, dS + \iint_S Q\hat{j} \cdot \hat{n} \, dS + \iint_S R\hat{k} \cdot \hat{n} \, dS. \end{aligned}$$

So we need to prove:

- $\iint_S P\hat{i} \cdot \hat{n} \, dS = \iiint_E \frac{\partial P}{\partial x} \, dV$
- $\iint_S Q\hat{j} \cdot \hat{n} \, dS = \iiint_E \frac{\partial Q}{\partial y} \, dV$
- $\iint_S R\hat{k} \cdot \hat{n} \, dS = \iiint_E \frac{\partial R}{\partial z} \, dV \leftarrow \text{proving this one}$

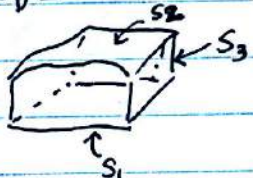
Since  $E$  is a type 1 region,  $E = \{(x, y, z) \mid (x, y) \in D, u_1(x, y) \leq z \leq u_2(x, y)\}$   
 $D = \text{projection of } E \text{ on } x\text{-}y \text{ plane.}$

From section 15.7 (Triple Integrals), we have:

$$\iiint_E \frac{\partial R}{\partial z} \, dV = \iint_D \left[ \int_{u_1(x, y)}^{u_2(x, y)} \frac{\partial R}{\partial z}(x, y, z) \, dz \right] \, dA \quad \text{because } \frac{\partial R}{\partial z} \text{ is func. of } dz \text{ 3 var.}$$

In the inner integral,  $x$  and  $y$  are treated like constants.  
 By FTC,  $\int_{u_1(x, y)}^{u_2(x, y)} \frac{\partial R}{\partial z}(x, y, z) \, dz = R(x, y, u_2(x, y)) - R(x, y, u_1(x, y))$ .

$$* \text{ So } \iiint_E \frac{\partial R}{\partial z} \, dV = \iint_D [R(x, y, u_2(x, y)) - R(x, y, u_1(x, y))] \, dA.$$

$\iint_S R\hat{k} \cdot \hat{n} \, dS:$    $S_1: z = u_1(x, y)$   
 $S_2: z = u_2(x, y)$   
 $S_3: \text{between } S_1 \text{ \& } S_2. \text{ "lateral" surface}$

$$\text{So } \iint_S R\hat{k} \cdot \hat{n} \, dS = \iint_{S_1} R\hat{k} \cdot \hat{n} \, dS + \iint_{S_2} R\hat{k} \cdot \hat{n} \, dS + \iint_{S_3} R\hat{k} \cdot \hat{n} \, dS.$$

On  $S_3$ ,  $\hat{n}$  is horizontal, so  $\hat{n} \cdot \hat{k} = 0$ .

On  $S_2$ ,  $\hat{n}$  points upward.

On  $S_1$ ,  $\hat{n}$  points downward.

Using  $\iint_S \vec{F} \cdot d\vec{S} = \iint_S \vec{F} \cdot \vec{n} \, dS = \iint_D (-P \frac{\partial g}{\partial x} - Q \frac{\partial g}{\partial y} + R) \, dA$   
 where  $\vec{F} = \langle 0, 0, R \rangle$  and  $u_2 = g$  we get:

$$\iint_{S_2} R \hat{k} \cdot \vec{n} \, dS = \iint_D R \, dA = \iint_D R(x, y, u_2(x, y)) \, dA$$

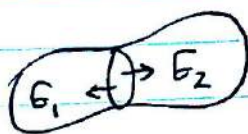
When  $u_1 = g$ , we get:

$$\iint_{S_1} R \hat{k} \cdot \vec{n} \, dS = \iint_D R \, dA = -\iint_D R(x, y, u_1(x, y)) \, dA.$$

Combining, we get:

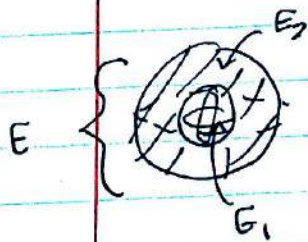
$$\iint_S R \cdot \hat{k} \cdot \vec{n} \, dS = \iint_D [R(x, y, u_2(x, y)) - R(x, y, u_1(x, y))] \, dA \checkmark \\ = \iiint_E \frac{\partial R}{\partial z} \, dV.$$

Extension to more general Regions:



Only the outer surface "survives" so

$$\sum_{i=1}^n \iiint_{E_i} \text{div } \vec{F} \, dV = \iiint_E \text{div } \vec{F} \, dV.$$



You can "add" regions.

$$E_2 = E - E_1 \\ \iint_{E_2} \vec{F} \cdot d\vec{S} = \iint_E \vec{F} \cdot d\vec{S} - \iint_{E_1} \vec{F} \cdot d\vec{S}.$$

Div and Flux: Sources and Sinks

Let  $\text{div } \vec{F}$  be continuous and  $P_0 = (x_0, y_0, z_0)$

Let  $B_a$  = ball w/ radius  $a$  and center  $P_0$ .

Let  $S_a$  = boundary of  $B_a$ .

$$\text{If } a \text{ is small, } \iint_{S_a} \vec{F} \cdot d\vec{S} = \iiint_{B_a} \text{div } \vec{F} \, dV \approx \iiint_{B_a} \text{div } \vec{F}(P_0) \, dV.$$

$$= \text{div } \vec{F}(P_0) \cdot V(B_a) = \text{div } \vec{F}(P_0) \cdot \frac{4}{3}\pi a^3.$$

gets better as  $a \rightarrow 0$ .

So,  $\text{div } \vec{F}(P_0)$  is the rate of outward flux per unit volume at  $P_0$ .

If  $\text{div } \vec{F}(P_0) > 0$ , the net flow is outward near  $P_0$ , and  $P_0$  is called a source.

If  $\text{div } \vec{F}(P_0) < 0$ , the net flow is inward near  $P_0$ , and  $P_0$  is called a sink.